Search for Gravitational-Wave Background from Cosmic Strings with PPTA DR2

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Introduction

- A pulsar timing array (PTA) offers a unique opportunity of detecting very low frequency gravitaional waves (GWs) from nHz to μ Hz, by regularly monitoring the time of arrivals (ToAs) of radio pulses from an array of highly stable millisecond pulsars in the Milky Way.
- Cosmic strings are linear topological defects that can either form in the early Universe from symmetry-breaking phase transitions at high energies or be the fundamental strings of superstring theory (or one-dimensional D-branes) stretched out to astrophysical lengths.

GWB from Cosmic Strings



- After their formation, the intersection between cosmic strings can lead to reconnections and form loops, which will then decay due to relativistic oscillation and emit gravitational waves.
- The gravitational-wave background (GWB) produced by cosmic strings could potentially detected by PTAs.

Dataset and Method

• The PPTA DR2 includes pulse ToAs from high-precision timing observations for 26 pulsars spanning about 15 years, with observations taken at a cadence of approximately three weeks.

We describe the GWB energy spectrum of cosmic strings in terms of the dimensionless tension, $G\mu$, and the reconnection probability, p. Even though p = 1 for classical strings, it can be less than 1 in the string-theory-inspired models or pure Yang-Mills theory. The dimensionless GW energy density parameter per logarithm frequency as the fraction of the critical energy density is

$$\Omega_{\rm gw}(f) = \frac{8\pi G f}{3H_0^2 p} \rho_{\rm gw}(t_0, f), \tag{1}$$

where t_0 is cosmic time today, and ρ_{gw} is the GW energy density per unit frequency that can be computed by

$$\rho_{\rm gw}(t,f) = G\mu^2 \sum_{n=1}^{\infty} C_n P_n,$$
(2)

 $C_n(f) = \int_0^{t_0} \frac{dt}{(1+z)^5} \frac{2n}{f^2} \mathbf{n}(l,t).$ (3)

Here, P_n is the radiation power spectrum of each loop, and n(l, t) is the density of loops per unit volume per unit range of loop length l existing at time t.

- The timing residuals $\delta \mathbf{t}$ for each pulsar are decomposed into
 - $\delta t = M \epsilon + \delta t_{\rm RN} + \delta t_{\rm DET} + \delta t_{\rm WN} + \delta t_{\rm CP}.$
 - $-M\epsilon$ accounts for the inaccuracies in the subtraction of timing model.
 - $-\delta t_{\rm RN}$ is the contribution from red noise, including achromatic spin noise, frequency-dependent dispersion measure noise, frequencydependent chromatic noise, achromatic band noise and system/group noise.
 - $-\delta t_{\rm DET}$ is deterministic noise, including chromatic exponential dips, extreme scattering events, and annual dispersion measure variations.
 - $\delta t_{\rm WN}$ is white noise, including a scale parameter on the TOA uncertainties (EFAC), an added variance (EQUAD), and a per-epoch variance (ECORR) for each backend/receiver system.

Results

with



- We first consider a model in which both the string tension $G\mu$ and the reconnection probability p are free parameters.
- The Bayes factor of the model including both the UCP and cosmic string signal versus the model including only the UCP is $\mathcal{BF}_{\text{UCP}}^{\text{UCP}+\text{CS}} = 0.591 \pm 0.008 < 3$, indicating no evidence for a GWB signal produced by the cosmic string in the PPTA DR2.

 $-\delta t_{\rm CP}$ is the contribution from the common-spectrum process (such as a GWB).

References

- [1] Z. C. Chen, Y. M. Wu and Q. G. Huang, "Search for the Gravitational-Wave Background from Cosmic Strings with the Parkes Pulsar Timing Array Second Data Release," arXiv:2205.07194
- [2] M. Kerr, D. J. Reardon, G. Hobbs, et al. "The Parkes Pulsar Timing Array project: second data release," Publ. Astron. Soc. Austral. 37, e020 (2020)



- We also consider models in which p is fixed while $G\mu$ is allowed to vary.
- For all of the values $p \in [10^{-3}, 1]$, we have $\mathcal{BF}_{UCP}^{UCP+CS} \leq 3$, confirming that there is no evidence for a GWB produced by cosmic strings in the PPTA DR2.
- We place 95% upper limit on cosmic string tension $G\mu$ as a function of p. For p = 1, the upper bound is $G\mu \lesssim 5.1 \times 10^{-10}$.